



AD FALCON API Manual

MIT-S1 Model (Pestana & Whittle, 1999)

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1 MIT-S1 Model (Pestana & Whittle, 1999)

MIT-S1 is a critical-state, bounding-surface constitutive model for soils. In FALCON it is available through the `MITS1Model` UMAT. This page summarizes the governing equations, input parameters, custom state variables, and the single-point benchmark cases included with the manual.

1.1 Syntax

This model is configured in `% Materials` as a user-defined mechanical material. Use `@UMAT:` with category `Mechanical` and pass the parameters as `name=value` pairs.

Example:

```
@UMAT: path/to/MITS1Model.cpp path/to/MITS1Model.hpp Mechanical \
  Pa=100 rho_c=0.25 p_ref=100 theta=0.20 \
  Cb=750 K0NC=0.50 nu0=0.25 omega=1.25 omega_s=4.0 \
  phi_cs=32 phi_mr=30 p=2.0 m=0.8 psi=30 \
  D=0 r=1 h=0 enableOCMapping=1 \
  P_min=1e-9 FTOL=1e-4 stressRelTol=1e-5 maxSubsteps=500 \
  CustomVariable=Alpha,b_xx,b_yy,b_zz,b_zy,b_zx,b_xy,alpha0_star,
alpha0_i,p_srp,srp_has_reversal,eta_srp_xx,eta_srp_yy,eta_srp_zz,
eta_srp_zy,eta_srp_zx,eta_srp_xy,eps_srp_xx,eps_srp_yy,eps_srp_zz,
eps_srp_zy,eps_srp_zx,eps_srp_xy
```

For readability, this example is wrapped across multiple lines; in input files you should write the full `@UMAT:` directive on a single line.

Example inputs used on this page:

- Generic `% Materials` snippet: [materials_mits1.txt](#)
- Pestana & Whittle (1999) Fig. 13-style (K_0) undrained clay case: [pestana1999_fig_13_ck0_strength_matching.txt](#)
- Pestana & Whittle (1999) Fig. 15-style (K_0) undrained clay case: [pestana1999_fig_15_ck0_psi30.txt](#)

1.2 Material parameters

Table 1. Core MIT-S1 material parameters

Symbol	Keyword in input	Required	Default	Role in the model
p_a	Pa	✓	none	Reference pressure used to normalise stiffness and stress levels in Eqs. (1), (15), and (19).
ρ_c	rho_c	✓	none	LCC compressibility; controls the slope of the limiting compression curve in Eq. (2) and enters isotropic hardening and flow through Eqs. (9), (13), and (25b).
p'_{ref}	p_ref	✓	none	Fixes the LCC position in Eq. (2), so it sets where isotropic compression and δ_b calculations are anchored.
θ	theta	✓	none	Transition exponent governing how quickly the response approaches the LCC in Eqs. (1), (8b), (9), and (25b).
C_b	Cb	✓	none	Elastic stiffness parameter controlling K_{max} and ρ_s in Eqs. (15) and (19).

Symbol	Keyword in input	Required	Default	Role in the model
K_0^{NC}	K0NC	✓	none	Defines the normally consolidated stress ratio and enters the anisotropy/flow scalars in Eqs. (8c) and (13).
ν_0	nu0	✓	none	Poisson ratio at stress reversal; sets the elastic ratio $2G_{\max}/K_{\max}$ in Eq. (15b) and the evolving ν in Eq. (18).
ω	omega	✓	none	Controls shear/Poisson degradation away from the SRP through Eqs. (16), (18), and (20).
ω_s	omega_s	✓	none	Controls small-strain shear degradation and the swelling slope in Eqs. (19b) and (20).
ϕ'_{cs}	phi_cs	✓	none	Critical-state friction angle; defines the failure surface in Eq. (3), the LCC coupling constant α_0^2 in Eq. (10), and the reference branch in Eq. (5).

Symbol	Keyword in input	Required	Default	Role in the model
ϕ'_{mr}	phi_mr	✓	none	Reference peak friction angle at $e = 1$; controls the aperture scalar through Eqs. (4c) and (5).
n_p	p	✓	none	Void-ratio exponent in Eq. (5); governs how strongly density changes the peak friction angle.
m	m	✓	none	Bounding-surface slenderness parameter; appears directly in the surface shape, flow, and mapping laws in Eqs. (4a), (7), (12a), and (24a).
ψ	psi	✓	none	Kinematic hardening rate of the anisotropy tensor b in Eq. (7); it controls how fast the surface translates.
D	D	×	0	Hysteretic swelling contribution after stress reversal in Eq. (19b).
r	r	×	1	Exponent controlling the decay of the hysteretic term $D(1 - \mu^r)$ in Eq. (19b).

Symbol	Keyword in input	Required	Default	Role in the model
h	h	×	0	Overconsolidated mapping modulus scale in Eq. (25b).

Table 2. Numerical and advanced controls

Keyword in input	Required	Default	Role
P_min	✓	none	Positive floor for p' and α' to avoid singular stress-ratio updates.
FTOL	✓	none	Admissibility tolerance for the yield function and local correction steps.
stressRelTol	×	1e-5	Relative stress target used by the adaptive substepping error estimator.
maxSubsteps	×	500	Maximum number of internal substeps per increment.
enableOC Mapping	×	1	Enables the interior bounding-surface mapping of Eqs. (22)–(25).
Alpha_max	×	1e12	Safety cap on α' during numerical integration.
hardeningRate Cap	×	50	Caps logarithmic hardening rates as a numerical safeguard.
alphaOverPMax	×	0	Optional cap on α'/p' ; 0 disables the cap.
alphaCapScale	×	0	Optional LCC-based cap on α' growth; 0 disables it.
initOverride	×	0	Rebuilds internal state from the current stress/void-ratio state during conditioning.
initOCR	×	1	OCR-style scale used when <code>initOverride=1</code> .

Keyword in input	Required	Default	Role
initBScale	×	1	Scale factor for rebuilding b from the current stress ratio.
initAlphaMode	×	1	Chooses how conditioning reconstructs α' from the current state.
initEnforce Admissible	×	1	Enlarges α' during conditioning if the current state is inadmissible.

1.3 Custom state variables

Declare custom state variables using CustomVariable= so they are stored for restart/output.

Table 3. MIT-S1 custom variables

Keyword(s)	State variable	Required	Role
Alpha	α'	✓	Current bounding-surface size used in Eqs. (4a), (6b), (6c), (8b), and (12), and updated by Eqs. (9) to (11).
b_xx, b_yy, b_zz, b_zy, b_zx, b_xy	b	✓	Components of the anisotropy tensor that translate the surface in Eq. (4a) and evolve through Eq. (7).
alpha0_star	α_0^*	✓	Current constant- η image-point size from Eq. (24a), used in the mapping factor g_1 in Eq. (24).

Keyword(s)	State variable	Required	Role
alpha0_i	α_{0i}	✓	First-yield loading-surface size used with α_0^* in Eq. (24) during overconsolidated mapping.
p_srp	p'_{srp}	✓	Mean effective stress at the stored stress-reversal point, used in Eqs. (16), (17), and (19b).
srp_has_reversal	SRP flag	✓	Indicates whether the stress-reversal-point history has been initialized for the hysteretic elasticity update.
eta_srp_xx, eta_srp_yy, eta_srp_zz, eta_srp_zy, eta_srp_zx, eta_srp_xy	η_{srp}	✓	Stress-ratio tensor at the stress-reversal point, used to evaluate m_4 in Eq. (17).
eps_srp_xx, eps_srp_yy, eps_srp_zz, eps_srp_zy, eps_srp_zx, eps_srp_xy	ϵ_{srp}	✓	Strain tensor at the stress-reversal point, used by the loading-unloading detector in Eq. (21).

1.4 Stress invariants and notation

Stress invariants are written in effective-stress form:

$$p' = \frac{1}{3} \text{tr}(\boldsymbol{\sigma}'), \quad \boldsymbol{s} = \boldsymbol{\sigma}' - p' \boldsymbol{I}, \quad \boldsymbol{\eta} = \frac{\boldsymbol{s}}{p'}. \quad (\text{o})$$

The third invariant of the stress-ratio tensor is

$$J_3^\eta = \det(\boldsymbol{\eta}) = \frac{\det(\boldsymbol{s})}{(p')^3}. \quad (\text{oa})$$

The internal variables are the void ratio e , the scalar bounding-surface size α' , and the deviatoric anisotropy tensor \mathbf{b} .

Plastic flow is written as

$$d\varepsilon^p = d\lambda \mathbf{P}, \quad d\varepsilon_v^p = d\lambda P_p, \quad d\varepsilon_d^p = d\lambda \mathbf{P}_s. \quad (ob)$$

1.5 LCC Compression Model (Eqs. 1–2)

For hydrostatic first loading:

$$d\varepsilon_v = \frac{e}{1+e} \left[\frac{\delta^\theta}{C_b \left(\frac{p'}{p_a}\right)^{1/3}} + \frac{\rho_c}{\left(\frac{p'}{p_a}\right)} (1 - \delta^\theta) \right] \frac{dp'}{p_a}, \quad (1a)$$

$$d\varepsilon_v^e = \frac{e}{1+e} \frac{1}{C_b \left(\frac{p'}{p_a}\right)^{1/3}} \frac{dp'}{p_a}. \quad (1b)$$

The LCC distance measure is

$$\delta = 1 - \frac{p'}{p'_b}, \quad p'_b = p'_{\text{ref}} \left(\frac{1}{e}\right)^{1/\rho_c}. \quad (2)$$

1.6 Critical-State Failure and Bounding Surface (Eqs. 3–5)

The Matsuoka-Nakai critical-state failure function is

$$h_f(\boldsymbol{\eta}) = k^2 - \boldsymbol{\eta} : \boldsymbol{\eta} = 0, \quad (3a)$$

with

$$k^2 = k_a^2 + \left(3 - \frac{k_a^2}{2}\right) J_3^\eta, \quad k_a^2 = \frac{8 \sin^2 \phi'_{cs}}{3 + \sin^2 \phi'_{cs}}. \quad (3b)$$

The distorted lemniscate bounding surface is

$$f = (p')^2 \left[(\boldsymbol{\eta} - \mathbf{b}) : (\boldsymbol{\eta} - \mathbf{b}) - \zeta^2 \left(1 - \left(\frac{p'}{\alpha'}\right)^m\right) \right] = 0, \quad (4a)$$

$$\zeta^2 = c^2 + \mathbf{b} : \mathbf{b} - 2 \boldsymbol{\eta} : \mathbf{b}, \quad (4b)$$

$$c^2 = c_a^2 + \left(3 - \frac{c_a^2}{2}\right) J_3^\eta, \quad c_a^2 = \frac{8 \sin^2 \phi'_m}{3 + \sin^2 \phi'_m}. \quad (4c)$$

The density-dependent peak friction angle is

$$\frac{1}{\tan \varphi'_m} = \frac{1}{\tan \left(45^\circ + \frac{\varphi'_{cs}}{2} \right)} + \left[\frac{1}{\tan \varphi'_{mr}} - \frac{1}{\tan \left(45^\circ + \frac{\varphi'_{cs}}{2} \right)} \right] e^{n_p}. \quad (5)$$

Here e^{n_p} means void ratio raised to the power n_p , not an exponential.

1.7 Yield Gradient and Hardening (Eqs. 6–11)

Define

$$\mathbf{Q} = \frac{\partial f}{\partial \boldsymbol{\sigma}'} \equiv (Q_p, \mathbf{Q}_s). \quad (6a)$$

With $g = (p'/\alpha')^m$, the components used in the UMAT are

$$Q_p = p' \left[(m\zeta^2 + 2 \boldsymbol{\eta} : \mathbf{b})g - 2 \boldsymbol{\eta} : \boldsymbol{\eta} + \left(9 - \frac{3c_a^2}{2} \right) (1-g)J_3^\eta \right], \quad (6b)$$

$$\mathbf{Q}_s = p' \left[2(\boldsymbol{\eta} - g\mathbf{b}) - \left(3 - \frac{c_a^2}{2} \right) (1-g) \frac{\partial J_3^\eta}{\partial \boldsymbol{\eta}} \right]. \quad (6c)$$

Kinematic hardening of the anisotropy tensor follows

$$d\mathbf{b} = \psi \frac{1+e}{e \alpha'} \left[\frac{r_x}{m} \langle Q_p d\varepsilon_v^p \rangle + r_y |\mathbf{Q}_s : d\varepsilon_d^p| \right] (\boldsymbol{\eta} - \mathbf{b}), \quad (7)$$

with Macaulay brackets $\langle x \rangle = \max(x, 0)$ and

$$r_x = \frac{k^2 + \mathbf{b} : \mathbf{b} - 2 \boldsymbol{\eta} : \mathbf{b}}{k_a^2}, \quad (8a)$$

$$r_y = (d^2 + \mathbf{b} : \mathbf{b} - 2 \boldsymbol{\eta} : \mathbf{b}) \left[1 + \left(\frac{\alpha'}{p'} - 1 \right) \delta_b^\theta \right], \quad (8b)$$

$$d^2 = d_a^2 + \left(3 - \frac{d_a^2}{2} \right) J_3^\eta, \quad d_a^2 = \frac{2(1 - K_0^{NC})^2}{1 + K_0^{NC} + (K_0^{NC})^2}. \quad (8c)$$

The isotropic hardening law for α' is

$$\frac{d\alpha'}{\alpha'} = \frac{1+e}{e(\rho_c - \rho_s)(1 - \delta_b)^\theta} \left[d\varepsilon_v^p + \delta_b^\theta \left(\frac{\mathbf{Q}_s : d\varepsilon_d^p}{p'} \right) \right] - \frac{2 \mathbf{b} : d\mathbf{b}}{\alpha_0^2 + \mathbf{b} : \mathbf{b}}, \quad (9)$$

where

$$\delta_b = 1 - \frac{\alpha'}{p'_b} \left(1 + \frac{\mathbf{b} : \mathbf{b}}{\alpha_0^2} \right), \quad \alpha_0^2 = 24 \left(\frac{\sin \varphi'_{cs}}{3 - \sin \varphi'_{cs}} \right)^2. \quad (10)$$

In the LCC regime ($\delta_b = 0$),

$$\frac{d\alpha'}{\alpha'} = \frac{1+e}{e(\rho_c - \rho_s)} d\varepsilon_v^p - \frac{2\mathbf{b} : d\mathbf{b}}{\alpha_0^2 + \mathbf{b} : \mathbf{b}}. \quad (11)$$

1.8 Non-Associated Flow and Consistency (Eqs. 12–14)

The volumetric flow component is

$$P_p = \begin{cases} (k^2 - \boldsymbol{\eta} : \boldsymbol{\eta}) \frac{p'}{\alpha'} (1 - \delta_b)^m, & \boldsymbol{\eta} : \boldsymbol{\eta} \leq k^2, \\ (k^2 - \boldsymbol{\eta} : \boldsymbol{\eta}) \frac{p'}{\alpha'}, & \boldsymbol{\eta} : \boldsymbol{\eta} > k^2, \end{cases} \quad (12a)$$

and the deviatoric flow component is

$$\mathbf{P}_s = x P_p \boldsymbol{\eta} + \frac{\zeta^2 \|\boldsymbol{\eta}\|}{\alpha'} \mathbf{Q}_s. \quad (12b)$$

The coefficient imposing the K_0^{NC} constraint is

$$x = \left(\frac{\rho_c}{\rho_c - \rho_s} \right) \left[\frac{1 + 2K_0^{NC}}{3(1 - K_0^{NC})} - \frac{K}{2G} \frac{\rho_s}{\rho_c} \right]. \quad (13)$$

Consistency gives

$$d\lambda H = -\frac{\partial f}{\partial \alpha'} d\alpha' - \frac{\partial f}{\partial \mathbf{b}} : d\mathbf{b}, \quad (14a)$$

$$d\lambda = \frac{K Q_p d\varepsilon_v + 2G \mathbf{Q}_s : d\varepsilon_d + \frac{\partial f}{\partial e} (1+e) d\varepsilon_v}{H + K Q_p P_p + 2G \mathbf{Q}_s : \mathbf{P}_s}. \quad (14b)$$

1.9 Hysteretic Elasticity (Eqs. 15–21)

At the stress-reversal point,

$$\frac{K_{\max}}{p_a} = C_b \left(\frac{1+e}{e} \right) \left(\frac{p'}{p_a} \right)^{1/3} \left(1 + \frac{K_{\max}}{2G_{\max}} \boldsymbol{\eta} : \boldsymbol{\eta} \right)^{1/6}, \quad (15a)$$

$$\frac{2G_{\max}}{K_{\max}} = 3 \left(\frac{1 - 2\nu_0}{1 + \nu_0} \right). \quad (15b)$$

The shear/bulk ratio degrades with distance from the SRP:

$$\frac{(2G/K)}{(2G_{\max}/K_{\max})} = \begin{cases} \frac{1}{1 + \omega m_4}, & p' < p'_{\text{srp}}, \\ \frac{1}{1 + \omega \mu m_4}, & p' \geq p'_{\text{srp}}, \end{cases} \quad (16)$$

with

$$\mu = \begin{cases} p'/p'_{\text{srp}}, & p' < p'_{\text{srp}}, \\ p'_{\text{srp}}/p', & p' \geq p'_{\text{srp}}, \end{cases} \quad m_4 = [(\eta - \eta_{\text{srp}}) : (\eta - \eta_{\text{srp}})]^{1/2}. \quad (17)$$

The equivalent Poisson ratio is

$$\nu = \nu_0 + \frac{\frac{1}{3}\omega m_4(1 + \nu_0)}{1 + \frac{2}{3}\omega m_4(1 + \nu_0)}, \quad \nu_0 \leq \nu < 0.5. \quad (18)$$

The tangent bulk modulus and swelling slope are

$$\frac{K}{p_a} = \frac{1 + e}{e \rho_s} \frac{p'}{p_a}, \quad (19a)$$

$$\rho_s = D(1 - \mu^r) + \frac{1 + \omega_s m_4}{C_b} \left(1 + \frac{K_{\max}}{2G_{\max}} \eta : \eta\right)^{1/6} \left(\frac{p'}{p_a}\right)^{2/3}. \quad (19b)$$

and the small-strain shear reduction is

$$\frac{G}{G_{\max}} = \frac{1}{(1 + \omega m_4)(1 + \omega_s m_4)}. \quad (20)$$

The SRP is identified from the strain increment relative to the previous reversal state:

$$s ds = \begin{cases} \Delta \varepsilon_v d\varepsilon_v, & d\varepsilon_v \neq 0, \\ \Delta \varepsilon_d : d\varepsilon_d, & d\varepsilon_v = 0, \end{cases} \quad \begin{cases} ds > 0 & \text{loading,} \\ ds < 0 & \text{unloading (new SRP).} \end{cases} \quad (21)$$

1.10 Bounding-Surface Mapping for Overconsolidated States (Eqs. 22–25)

Loading is detected from the image-point gradient:

$$KQ_p^I d\varepsilon_v + 2GQ_s^I : d\varepsilon_d \begin{cases} \geq 0 & \text{loading,} \\ < 0 & \text{unloading.} \end{cases} \quad (22)$$

When mapping is active,

$$\mathbf{P} = (1 - g_1)\mathbf{P}^I + g_1\mathbf{P}_0, \quad (23a)$$

$$H = \langle H^I \rangle + H_0 \frac{g_1}{1 - g_1} \left(1 - \frac{\eta : \eta}{c^2}\right)^{1/2}, \quad (23b)$$

$$g_1 = \frac{\alpha' - \alpha_0^*}{\alpha' - \alpha_{0i}}, \quad 0 \leq g_1 \leq 1, \quad (24)$$

with first-yield quantities

$$P_{0p} = -2\|\boldsymbol{\eta} - \mathbf{b}\| \left\| \frac{\mathbf{s}^I}{\alpha'} \right\|, \quad P_{0s} = P_s^I, \quad (25a)$$

$$H_0 = \left(\frac{\rho_s^I}{\rho_c - \rho_s^I} \right) \frac{h}{1 - \delta_b^\theta} K_{\max}^I \|\mathbf{Q}^I\| \|\mathbf{P}^I\|. \quad (25b)$$

The loading-surface size stored in `alpha0_star` is evaluated from the constant- η image-point relation

$$\alpha_0^* = \frac{p'}{(g^*)^{1/m}}, \quad g^* = 1 - \frac{(\boldsymbol{\eta} - \mathbf{b}) : (\boldsymbol{\eta} - \mathbf{b})}{\zeta^2}. \quad (24a)$$

1.11 Implementation Notes

- The constitutive equations above define the model. In the public FALCON interface, MIT-S1 uses explicit adaptive substepping for local stress integration.
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1.12 Single-Point Validation

This page includes two K_0 -undrained single-point benchmark reproductions based on Pestana & Whittle (1999). The figures compare the data reported in that paper with the corresponding FALCON MIT-S1 responses.

1.12.1 Prescribed K_0 consolidation effect on undrained clay response

The first reproduction follows the clay benchmark reported in Pestana & Whittle (1999) Figure 13. The figure compares the paper data with the FALCON MIT-S1 response for the same parameter sets.

Use [pestana1999_fig13_ck0_strength_matching.txt](#) as the starting deck and vary `phi_mr`, `m`, and the sign of `dEpsAxial` to recover the different compression/extension branches.

1.12.2 Anisotropy-rate effect in K_0 -consolidated undrained clay

The second reproduction follows Pestana & Whittle (1999) Figure 15, which isolates the effect of the anisotropy-rate parameter ψ for K_0 -consolidated clay. The figure compares the paper data with the FALCON MIT-S1 response.

MIT-S1 reproduction of Pestana & Whittle (1999) Figure 13

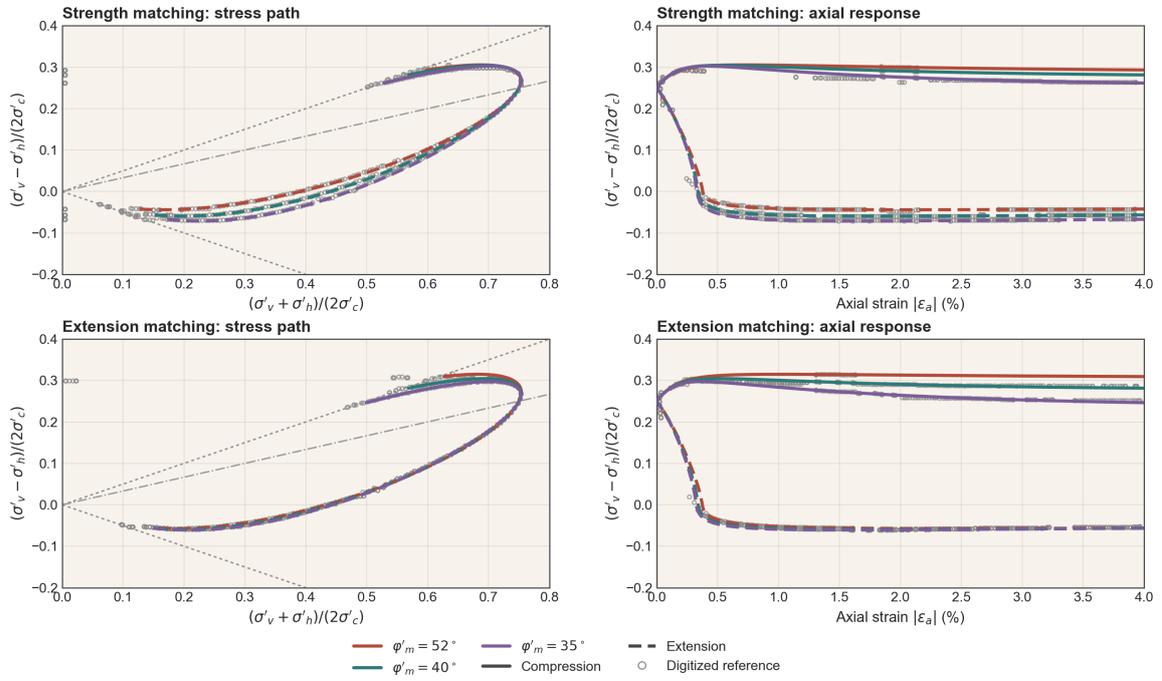


Figure 1: MIT-S1 alternate reproduction of Pestana and Whittle 1999 Figure 13

MIT-S1 reproduction of Pestana & Whittle (1999) Figure 15

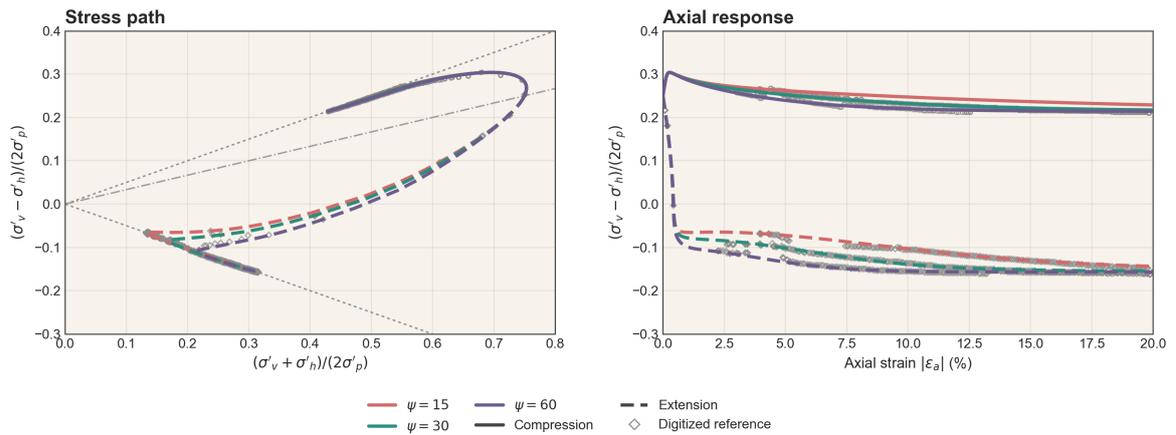


Figure 2: MIT-S1 alternate reproduction of Pestana and Whittle 1999 Figure 15

The sample deck [pestana1999_fig15_ck0_psi30.txt](#) reproduces the middle $\psi = 30$ case; sweep $\psi = 15, 30, 60$ to rebuild the full family.

1.13 References

- Pestana, J. M., & Whittle, A. J. (1999). *Formulation of a unified constitutive model for clays and sands*. *International Journal for Numerical and Analytical Methods in Geomechanics*, 23, 1215–1243.

